

## A gravitational effective action on a finite triangulation

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ABSTRACT: We construct a function of the edge-lengths of a triangulated surface whose variation under a rescaling of all the edges that meet at a vertex is the defect angle at that vertex. We interpret this function as a gravitational effective action on the triangulation, and the variation as a trace anomaly.

KEYWORDS: 2D Gravity, Lattice Models of Gravity.

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**1. Introduction**

Anomalies are generally regarded as arising from the infinite numbers of degrees of freedom in continuous systems. For example, the action  $S$  of a scalar field  $\Phi$  coupled to a gravitational background on a two dimensional surface  $\Sigma$ ,

$$S = \frac{1}{2} \int_{\Sigma} d^2x \sqrt{\det(g_{pq})} g^{mn} \partial_m \Phi \partial_n \Phi , \tag{1.1}$$

has a classical symmetry under rescalings of the metric on  $\Sigma$ :  $g_{mn} \rightarrow \lambda(x)g_{mn}$ . This implies

$$g_{mn} \frac{\partial S}{\partial g_{mn}} = 0 . \tag{1.2}$$

Upon quantization, this symmetry is anomalous; that is, if one defines the quantum effective action  $\Gamma$  as

$$e^{-\Gamma[g]} = \int [d\Phi] e^{-S[\Phi,g]} , \tag{1.3}$$

then one finds [1]

$$g_{mn} \frac{\partial \Gamma}{\partial g_{mn}} = -\frac{1}{24\pi} \sqrt{\det(g_{mn})} R(g(x)) , \tag{1.4}$$

where  $R(g(x))$  is the scalar curvature<sup>1</sup> of the two dimensional metric  $g_{mn}$ . If one integrates this over the surface, one finds

$$\int d^2x g_{mn} \frac{\partial \Gamma}{\partial g_{mn}} = -\frac{1}{24\pi} \int d^2x \sqrt{\det(g_{mn})} R(g(x)) = \frac{1}{6} \chi , \tag{1.5}$$

where  $\chi = 2(1 - g)$  is the Euler character of the two dimensional surface.

<sup>1</sup>We use the convention that the scalar curvature is minus twice the Gaussian curvature, and hence is negative on the sphere [1].

Both the Euler character  $\chi$  and the density  $\sqrt{g}R$  have sensible analogs on a triangulation of a surface:  $\chi = V - E + F$ , where  $V, E, F$  are the numbers of vertices, edges, and faces in the triangulation, respectively. Since  $\chi = \frac{1}{2\pi} \sum_{i \in \{V\}} \epsilon_i$  where  $i$  runs over all vertices and  $\epsilon_i$  is the defect at the  $i$ 'th vertex, one can identify the defect  $\epsilon_i = -\frac{1}{2}\sqrt{g}R_i$  with the curvature at the vertex [2]. Thus one has the correspondence:

$$-\frac{1}{4\pi} \int d^2x \sqrt{g}R = \chi \leftrightarrow \frac{1}{2\pi} \sum_{i \in \{V\}} \epsilon_i = V - E + F . \tag{1.6}$$

It is natural to ask if this correspondence can be extended to the anomaly (1.4): can one find an analog of the effective action  $\Gamma$  on a triangulation? That is, can one find a function  $\Gamma(l_{ij})$  of the edge lengths  $l_{ij}$  such that

$$\sum_{j \in \langle ij \rangle} l_{ij} \frac{\partial \Gamma}{\partial l_{ij}} = \frac{1}{12\pi} \epsilon_i \tag{1.7}$$

for all vertices  $i$  (the sum is over all edges with one end at  $i$ ).

In this note, we show that this is possible, and explicitly construct  $\Gamma$ . Our result is

$$\Gamma = \frac{1}{12\pi} \left[ \sum_{\angle_{ijk}} \int_{\frac{\pi}{2}}^{\alpha_{ijk}} \left(y - \frac{\pi}{3}\right) \cot(y) dy + \sum_{\langle ij \rangle} 2k_{ij}\pi \ln\left(\frac{l_{ij}}{l_0}\right) \right], \tag{1.8}$$

where the first sum is over all internal angles  $\alpha_{ijk}$ , the second sum is over all edges  $\langle ij \rangle$  with lengths  $l_{ij}$  (the explicit factor of two arises because every edge is shared by two triangles),  $l_0$  is a scale that we set to equal to one from now on, and the  $k_{ij}$  are constants associated to the edges that satisfy

$$\sum_{j \in \langle ij \rangle} k_{ij} = 1 - \frac{n_i}{6}, \quad n_i = \sum_{j \in \langle ij \rangle} 1 \tag{1.9}$$

at every vertex  $i$ ; here  $n_i$  is the number of neighbors of the  $i$ 'th vertex. Note that the conditions (1.9) do not in general determine the constants  $k_{ij}$  uniquely; one could add a subsidiary condition, e.g., that  $\sum k_{ij}^2$  is minimized, to remove this ambiguity. Note also that the total Euler character, which comes from a *uniform* scaling of all lengths and thus does not change the angles  $\alpha_{ijk}$ , comes entirely from the last term, i.e., from  $\partial\Gamma/\partial l_0$ .

The strategy that we use to find this solution is as follows: we first consider the simplest case, a triangulation of the sphere with three vertices, three edges, and two faces, and prove the integrability conditions needed for  $\Gamma$  to exist are satisfied. We then find  $\Gamma$  for this case and show that it immediately generalizes to all triangulations with a certain homogeneity property, and finally generalize  $\Gamma$  to an arbitrary triangulation.

It would be interesting to complete the correspondence, and find a way to compute the result (1.8) as the anomalous effective action corresponding to a discrete analog of, e.g., the scalar action (1.1); a promising approach might be the work of S. Wilson on triangulated manifolds [3].

## 2. Integrability

We begin with a triangulation of the sphere with three vertices, three edges, and two (identical) faces (a triangular “pillow”); we label the edges by their lengths  $a, b, c$  and the opposite internal angles of the triangles by  $\alpha, \beta$  and  $\gamma$ , respectively. We also abbreviate

$$a \frac{\partial \Gamma}{\partial a} \equiv D_a(\Gamma) , \text{ etc.}, \tag{2.1}$$

and, for simplicity, drop an overall factor of  $1/(12\pi)$  in  $\Gamma$ . The defect at the vertex  $\alpha$  in this case is just  $2(\pi - \alpha)$ , etc. Using this notation, the equations that we want  $\Gamma$  to satisfy in the triangle are:

$$\begin{aligned} D_a(\Gamma) + D_b(\Gamma) &= 2(\pi - \gamma) , \\ D_a(\Gamma) + D_c(\Gamma) &= 2(\pi - \beta) , \\ D_b(\Gamma) + D_c(\Gamma) &= 2(\pi - \alpha) . \end{aligned} \tag{2.2}$$

If we add the first two equations, and subtract the third, we get  $D_a(\Gamma) = (\pi - \gamma - \beta + \alpha)$ . Since  $\gamma + \beta + \alpha$  is  $\pi$ , we get

$$D_a(\Gamma) = 2\alpha , \quad D_b(\Gamma) = 2\beta , \quad D_c(\Gamma) = 2\gamma . \tag{2.3}$$

The function  $\Gamma$  can exist only if

$$D_b(D_a(\Gamma)) = D_a(D_b(\Gamma)) . \tag{2.4}$$

Using

$$\alpha = \arccos \left( \frac{-a^2 + b^2 + c^2}{2bc} \right) , \quad \beta = \arccos \left( \frac{a^2 - b^2 + c^2}{2ac} \right) , \tag{2.5}$$

it is easy to see that (2.4) is satisfied. Thus the integrability conditions are satisfied for the triangle.

The tetrahedron and octahedron give results similar to those of the triangle. However, for a general triangulation, it is not easy to decouple the integrability conditions and reduce them to equations that may be checked straightforwardly. Instead, we construct  $\Gamma$  explicitly.

## 3. The effective action

Because the triangle is the simplest case, and can be related to any other system, we have examined it in detail. As noted above, in this case the defect at a vertex is directly related to the internal angle at that vertex; this suggests that in the general case, where the defect is related to the sum of the internal angles at a vertex,  $\Gamma$  should be just the sum of the  $\Gamma$  for each triangle. This is almost correct.

The basic strategy for the triangle was to rewrite the differential equations for  $\Gamma$  in terms of new variables: the angles  $\alpha, \beta$ , and the edge length  $c$  between them. One can integrate some of the equations and finally arrive at  $\Gamma$  on a single triangle  $\Delta$ :

$$\Gamma_{\Delta} = \sum_i \left[ \left( \alpha_i - \frac{\pi}{3} \right) \ln(\sin(\alpha_i)) - \int_{\frac{\pi}{2}}^{\alpha_i} \ln(\sin(y)) dy + k_i \pi \ln(a_i) \right] , \tag{3.1}$$

where  $\{a_1, a_2, a_3; \alpha_1, \alpha_2, \alpha_3\} = \{a, b, c; \alpha, \beta, \gamma\}$ , and  $k_i$  are constants associated to each edge. This can be simplified by integration by parts:

$$\Gamma_\Delta = \sum_i \left[ \int_{\frac{\pi}{2}}^{\alpha_i} \left( y - \frac{\pi}{3} \right) \cot(y) dy + k_i \pi \ln(a_i) \right]. \tag{3.2}$$

Note that all terms are expressed in terms of the internal angles  $\alpha_i$  except for the last term, which explicitly involves  $a_i$ . To prove that this is correct (and to determine  $k_i$ ), we differentiate  $\Gamma$ :

$$\begin{aligned} a_i \frac{\partial \Gamma}{\partial a_i} &\equiv D_{a_i} \Gamma = \sum_j \left[ \left( \alpha_j - \frac{\pi}{3} \right) \cot(\alpha_j) D_{a_i} \alpha_j \right] + k_i \pi \\ &= \sum_j \left[ - \left( \alpha_j - \frac{\pi}{3} \right) \frac{\cos(\alpha_j) D_{a_i} \cos(\alpha_j)}{1 - \cos^2(\alpha_j)} \right] + k_i \pi. \end{aligned} \tag{3.3}$$

Then the contribution of one triangle to the defect at vertex 1 is given by

$$\begin{aligned} (D_b + D_c) \Gamma &= - \left( \alpha - \frac{\pi}{3} \right) \frac{\cos(\alpha) (D_b + D_c) \cos(\alpha)}{1 - \cos^2(\alpha)} - \left( \beta - \frac{\pi}{3} \right) \frac{\cos(\beta) (D_b + D_c) \cos(\beta)}{1 - \cos^2(\beta)} \\ &\quad - \left( \gamma - \frac{\pi}{3} \right) \frac{\cos(\gamma) (D_b + D_c) \cos(\gamma)}{1 - \cos^2(\gamma)} + (k_b + k_c) \pi \end{aligned}$$

which can be explicitly evaluated using  $\gamma = \pi - \alpha - \beta$  and (2.5), and gives:

$$(D_b + D_c) \Gamma = \frac{\pi}{3} - \alpha + k_b \pi + k_c \pi. \tag{3.4}$$

This calculation works just as well for any triangle in a general triangulation to give  $\Gamma$  on any surface.

The constants  $k_i$  are assigned to every edge  $i$ . Clearly, in the case of the triangular pillow, there is a trivial solution to  $k_i = \frac{1}{3}$  for all  $i$  (recall that  $\pi - \alpha$  is half the defect at the vertex, but there are two triangles meeting at each vertex to sum over in this case). More generally, for any locally homogeneous triangulation in which all vertices have  $n$  nearest neighbors, we can choose

$$k_i = \frac{1}{n} - \frac{1}{6}. \tag{3.5}$$

However, for a general triangulation, the edge may connect vertices with different numbers of edges connecting to them. In this case, it is not necessarily trivial to find the appropriate values for all the  $k_i$ .

#### 4. A problem in graph theory

On a general triangulated surface, the we have the condition

$$\sum_{j \in \langle ij \rangle} k_{ij} = 1 - \frac{n_i}{6} \tag{4.1}$$

at every vertex, where  $n_i$  is the number of neighbors of the  $i$ 'th vertex. This means we want to find labels for the edges of a graph such that the sum at each vertex is the same. The equations (4.1) are a system  $V$  linear equations on the  $E$  variables  $k_{ij}$ , where  $V$  is

the total number of vertices and  $E$  is the total number of edges; note that  $E \geq V$  for all triangulations. We can rewrite this in terms of the  $V \times E$  dimensional matrix that describes the connections between vertices. This matrix has exactly two ones, which correspond to two vertices, in each column, which corresponds to the edge connecting the vertices.

We are happy to thank L. Motl [4] for the following proof that a solution to these equations always exists. The system of equations would not have a solution only if there were a linear combination of the rows that vanishes, that is, if there existed a vector that is perpendicular to all the columns. Because every column contains exactly two ones, it suffices to consider a submatrix that defines a triangle:

$$\begin{pmatrix} 1 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 1 \end{pmatrix}. \quad (4.2)$$

Since this matrix is nondegenerate, no nontrivial vector is annihilated by it. Since every edge sits on a triangle, no such vector can exist for the whole triangulation. Therefore, there must always be at least one way to label the edges.

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## References

- [1] A.M. Polyakov, *Quantum geometry of bosonic strings*, *Phys. Lett.* **B 103** (1981) 207.
- [2] T. Regge, *General relativity without coordinates*, *Nuovo Cim.* **19** (1961) 558.
- [3] S. Wilson, *Geometric structures on the cochains of a manifold*, 2005, [math.GT/0505227](#).
- [4] Luboš Motl, private communication.